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**Neutron and gamma-ray source-term characterization of AmBe sources
in Osaka University**

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Thermal / epi-thermal columns were constructed in the OKTAVIAN facility of Osaka University for basic studies of BNCT. For characterization of the columns, the spectrum and intensity of neutrons and the intensity of gamma-rays emitted from the AmBe sources used in the columns as a neutron source were measured. Neutrons were measured with the foil activation and multi-foil activation techniques, and gamma-rays were measured directly with an HpGe detector. The measured absolute intensity of neutrons is 2.4E6 n/sec. The measured neutron spectrum was fairly consistent with the previous result. The gamma-ray intensity was 1.8E6 photons/sec, which is a little smaller than that of neutrons. The measured data are used for normalization of various experiments carried out for BNCT.

Keywords: BNCT; source term; AmBe; multi-foil activation; HpGe detector; 4.44 MeV

1. Introduction

Boron neutron capture therapy (BNCT) is known to be a very effective cancer therapy. In Japan, clinical tests are underway using nuclear reactors. BNCT can kill only tumors by charged-particles produced by $^{10}\text{B}(n,\alpha)^7\text{Li}$ reaction. It is known that the effect to normal tissues could be suppressed substantially with BNCT.

However, there are problems recognized to be solved in BNCT. First, in Japan there exist only two nuclear reactor facilities available for BNCT. Accelerator based neutron sources (ABNS) are thus under investigation as a promising candidate for BNCT. At the moment, it is not yet realized because of various technical difficulties. The main reason is that the neutron intensity is weak and the energy is rather high. Patients should therefore be positioned very close to the target. As a result, the neutron spectrum would be strongly distorted and doses due to high energy neutrons and secondary gamma-rays become serious. The author's group is now developing a new device to characterize the neutron field with a position sensitive proportional counter for BNCT in ABNS as well as in nuclear reactors [1].

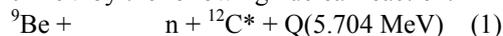
Also, a problem is known that it is difficult to know the exact BNCT dose in tumors. Currently, the irradiation time is determined with helps of simulation calculations and activation foils used during an irradiation. The key issue is that real-time monitoring of the BNCT dose is not possible. Our group is now

studying a new detection device with CdTe crystals for developing a SPECT system to know three-dimensional BNCT dose distribution in real time [2].

For these researches for BNCT, small and easy-to-use thermal and epi-thermal neutron fields are indispensable. We have carried out design studies to construct these neutron fields for last few years [3]. As a neutron source, AmBe was used. Last year we constructed two graphite columns and replaced some parts of them with other material to make thermal and epi-thermal neutron fields. However, the specification of the AmBe sources, i.e., neutron spectrum and intensity, is not known except Am radioactivity, because the AmBe sources are originally used as a startup neutron source for a reactor subcritical assembly of Osaka University in the past. The objectives of the study are to clarify the neutron source spectrum and intensity and in addition the intensity of simultaneously emitted gamma-rays of 4.44 MeV.

2. AmBe source

In the OKTAVIAN facility of Osaka University, there are four AmBe sources, which were used as a neutron source for operating a reactor subcritical assembly of Osaka University. The intensity of each is 46.3 GBq. The AmBe source generates neutrons up to around 10 MeV by the following nuclear reaction.



particles are emitted via decay of ^{241}Am .

According to Geiger [4], 70 neutrons per 10^6 particles is reported. In reality, the particle energy

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changes depending on the chemical form of Am and Be, and how to mix both of them. Since the energy change causes the neutron energy change, strictly speaking, it is necessary to measure the spectrum and intensity separately for each AmBe source stored in nuclear facilities.

The dimensions of the AmBe source are 3cm in diameter by 5cm in length. The source is double-sealed with stainless steel. Though the sources were quite old (fabricated in 1963), the intensity is almost the same as before because the half-life is as long as 400 years.

Also, from the excited residue nucleus, $^{12}\text{C}^*$, 4.44 MeV gamma-rays are emitted simultaneously. The AmBe can be used as a mono-energetic high energy gamma-ray source.

3. Neutron measurement

Neutron energy spectrum and intensity of the AmBe sources were measured with a multi activation foil technique. As a cross-check, the intensity was determined using a graphite thermal neutron column made with an AmBe source.

3.1. Measurement with multi activation foils

As is well known, neutron spectrum can be measured with a scintillation detector. However, in the case of AmBe, neutron and gamma-ray discrimination is needed, which is a little complicated technique. In the present study, the multi activation foil technique was employed for the neutron spectrum measurement, considering that the foil activation method is quite simple and convenient, and as a result it can easily be applied to many facilities having such kind of strong AmBe sources.

3.1.1 Selection of foils

The dynamic range of neutron energy for foils to be selected is up to around 11 MeV from the previous results. We examined threshold activation reactions up to 11 MeV. The criteria for the activation foil feature are set as follows: (1) Half-life of the foil should be longer than 5 minutes and shorter than 15 hours. (2) Evaluated cross section data for the activation foil should exist. (3) And, the number of counts more than several hundreds can be obtained by gamma-ray measurement for the period of half-life. As for the sample itself, the followings are taken into consideration: (1) Natural abundance of the nucleus to be activated should be more than 1 %. (2) Foils which cannot be obtained easily or are quite expensive should be avoided. (3) And, thin materials can be obtained by processing such that the correction factor due to thickness should be less than 20 %. From the above selection procedure, eight activation foils (nine reactions) were finally selected as shown in **Table 1**. **Figure 1** describes each activation reaction cross section. Fortunately, threshold energies for the nine reactions are enough scattering in order to reproduce the neutron energy.

Table 1. Selected foils.

Nucleus	Reaction	Effective threshold (MeV)
^{113}In	(n,n')	0.8
^{115}In	(n,n')	0.8
^{111}Cd	(n,n')	0.9
^{64}Zn	(n,p)	2.5
^{27}Al	(n,p)	4.2
^{56}Fe	(n,p)	6
^{24}Mg	(n,p)	6.4
^{27}Al	(n, α)	7
^{59}Co	(n, α)	7

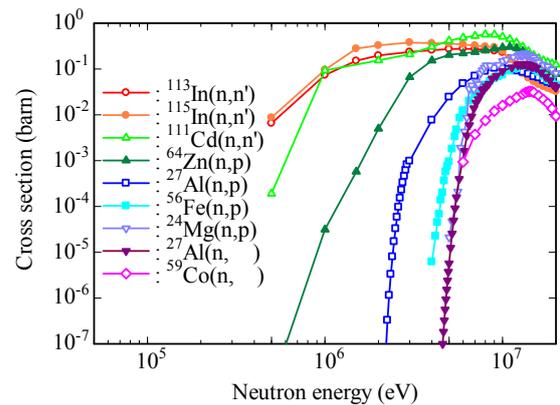


Figure 1. Cross sections of selected foils.

3.1.2 Irradiation and measurement

Experiments were conducted in “Large Experimental Room” of OKTAVIAN, which has a large experimental space of about $10 \times 10 \times 20 \text{ m}^3$. The AmBe sources stored in the OKTAVIAN facility are rather massive, i.e., 3cm in diameter and 5cm in length. However, since double-sealed structure is employed, the part where Am and Be materials are packed is fairly small. By a collimation experiment, it was confirmed that the source could be dealt with as a point source. Activation foils are thus located at 5 cm distant from the AmBe source. A foil is set on an expanded polyethylene behind the source. After irradiation, cooling for an appropriate time, gamma-rays emitted from activated foils were measured with an HpGe detector. The measuring time is basically fixed to be one half-life. Thousands of counts were obtained except ^{27}Mg and ^{56}Mn created from aluminum and cobalt samples, respectively.

3.1.3 Results

Measured reaction rates for eight activation foils were unfolded by a process based on the Bayes’ theorem with a response matrix for the foils evaluated from JENDL-Activation File [5]. The details of data processing are described elsewhere [6]. The result was compared with the previously measured spectrum [7] in **Figure 2**. It is noted that the cited reference spectrum is from ^{239}Pu -Be and was bunched so that the energy bin becomes the same as the present measurement. The agreement is acceptable. However, because the number of energy bins is a little small so that the expected fine

structure due to the nuclear levels of the target nucleus could not be confirmed. By integrating over energy, the source intensity can be estimated, which is 4.5×10^6 n/sec. This value is for the case with the source casing. According to Geiger, the measured value is 3.2×10^6 n/sec, and 3.8×10^6 n/sec is reported as a theoretically calculated value. The present result shows a little larger compared to these previous results. It should be however noted that the spectrum disagreement might also be caused by the difference of the chemical property of Am and Be.

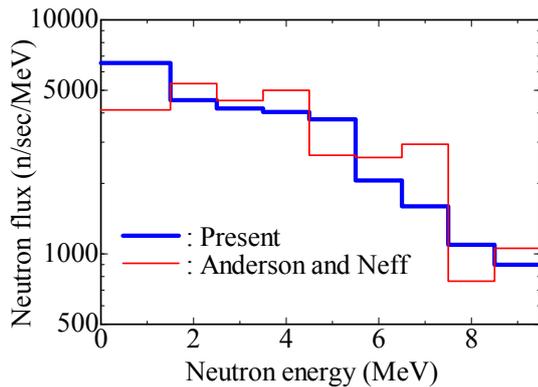


Figure 2. Measured neutron spectrum.

3.2. Cross-check by integral experiment

To check the disagreement with Geiger's value, an integral experiment with a graphite column was carried out. By this experiment the source intensity can be measured more accurately because one foil is enough in this kind of experiment and we can choose the one whose cross section accuracy is sufficiently high.

3.2.1 Design and construction of graphite column

At first, we designed a graphite column with MCNP5. An AmBe source was located in a cubic graphite column of $100 \times 100 \times 100$ cm³. Changing the distance between the source and the sample from 5 cm to 25 cm, the activation reaction rates of gold foils were calculated. As a result, the reaction rate shows its maximum at 5 cm and that of 15 cm shows a slight small value by 20%. Taking into account the fact that placement error of the AmBe source can increase the error of the present estimation, the foil was set at 15 cm.

3.2.2 Irradiation and simulation calculation

A graphite column was constructed in accordance with the design conducted in the previous section. An AmBe source and Au foil ($2 \times 2 \times 0.01$ cm³) were positioned and irradiation was performed for 3 days. After the irradiation, activity of the foil was measured with an HpGe detector. Also, transport calculations were carried out with a three dimensional model of the column including floor concrete by MCNP5. The source spectrum was assumed to be the one obtained in Sec. 3.1. With F4 tally, the reaction rate of $^{197}\text{Au}(n, \gamma)$ was estimated.

3.2.3 Results and discussion

By comparing experiment and calculation, the neutron source intensity was derived as 2.4×10^6 n/sec so that the reaction rate of $^{197}\text{Au}(n, \gamma)$ for both results became an equal value. This is the case without a casing. The obtained value is smaller than Geiger's and ours. Disagreement with Geiger is caused by the same reason mentioned in Sec. 3.1.3. The possible reason for the discrepancy within our results is thought to be uncertainty of the activation cross sections, i.e., all of the foils used in Sec. 3.1 have their evaluated activation reaction cross sections, however, some of them might not be validated sufficiently because they are not so frequently utilized. This may also be a cause of spectrum discrepancy discussed in Sec. 3.1.3. However, for the precise discussion for the spectrum, it is rather necessary to measure fine-energy-bin spectra by increasing the number of foils. The selection criteria in the present study were a little hard so that there are many foils which were abandoned. Contrary speaking, relaxing the criteria the number of foils can be increased. Also, direct measurements with a scintillation detector should be taken into account in the next step.

4. Gamma-ray measurement

From the AmBe source, 4.44 MeV gamma-rays are emitted simultaneously. Intensities of both neutrons and gamma-rays are thought to be more-or-less the same. In this section, the intensity of gamma-rays was measured. Not only when using the AmBe source as a gamma-ray source, but also even in case of using it as a neutron source, the gamma-ray intensity is quite important information.

4.1. Direct measurement with Hp-Ge detector

The measurement is quite simple. An AmBe source is placed in a large room and emitted gamma-rays are measured with an HpGe detector. The point is how to determine the detection efficiency for 4.44 MeV.

4.1.1 Detection efficiency

Detection efficiency for HpGe detector is easily fixed up to around 3 MeV because suitable standard gamma-ray sources are available. In the present case, it is needed to determine the efficiency at 4.44 MeV. For this purpose, the graphite column constructed in Sec. 3.2 was used. An iron sample was set in the column and gamma-rays emitted via $^{56}\text{Fe}(n, \gamma)$ reaction were measured. From an analysis with MCNP5, the optimum thickness of iron was 2cm and it was found that the gamma-ray emission efficiency became the best when the iron sample was positioned on the surface of the column and the AmBe source was set just behind the iron sample. The relative intensities of emitted gamma-rays were cited from the BNL database [8].

Figure 3 shows the obtained efficiency curve.

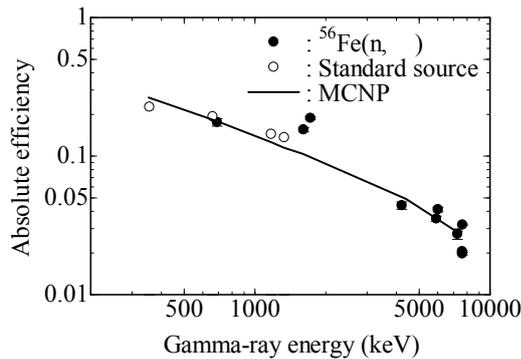


Figure 3. HpGe detector efficiency.

4.1.2. Spectrum and intensity measurement

Spectrum of gamma-rays emitted from the AmBe source was measured by the HpGe detector positioned at 3 m distant from the source. The measuring system was set close to the detector, but heavily shielded with 1m thick iron wall. The intensity of 4.44 MeV gamma-rays estimated from the net count of the photopeak was 1.8×10^6 photons/sec. This is a little smaller value than the neutron intensity of 2.4×10^6 n/sec for the case without a casing in Sec. 3.2.3.

4.2. Correlation with the neutron intensity

From the measured results, the intensity of gamma-rays shows 25% smaller than that of neutrons. The reason is discussed as follows: **Figure 4** shows the level scheme of ^{12}C . Now the α particle energy is 5.5 MeV. The level of compound nucleus of ^{13}C is ~ 11 MeV higher than the ground state of the exit channel of $^{12}\text{C} + n$. This means channels which should be taken into account are up to the five excited state of 10.8 MeV. Since gamma-rays of 4.44 MeV are emitted by a transition from the first excited state to the ground state of ^{12}C and only 4.44 MeV gamma-rays are measured, from other excited states a breakup reaction opens, which is a decay to three α . According to JENDL-AN-2005 [9], cross sections for possible decay channels of (α, n_{GS}) , $(\alpha, n_{1\text{st}})$, $(\alpha, n_{2\text{nd}})$ and $(\alpha, n_{\text{continuum}})$ are 50mb, 353mb, 39mb and 0b, respectively. For the channel of $(\alpha, n_{1\text{st}})$ emitting 4.44 MeV gamma-rays, the cross section is about 80 % out of the possible channels. In reality, α particles can induce nuclear reactions after slowing down. If the α particle energy is 2 MeV, the cross sections are going down and the ratio of the second reaction is about 75%. Though it is difficult to estimate the exact value, this is consistent with the measurement.

5. Conclusion

Spectrum and intensity of neutrons and intensity of gamma-rays emitted from the AmBe sources stored in the OKTAVIAN facility of Osaka University were measured. The neutron spectrum was fairly consistent with the previous result. However, precise discussion

was difficult because the number of energies was few. The absolute intensity was 2.4×10^6 n/sec, which was a little smaller than the value obtained before. The gamma-ray intensity showed 25 % smaller than that of neutron. In the next step, to estimate a finer-energy-bin neutron spectrum, we will carry out measurements with a scintillation detector, i.e., pulse height spectrum measurement, time-of-flight measurement and double scintillator measurement. The measured data will be used for various basic researches for BNCT in the future.

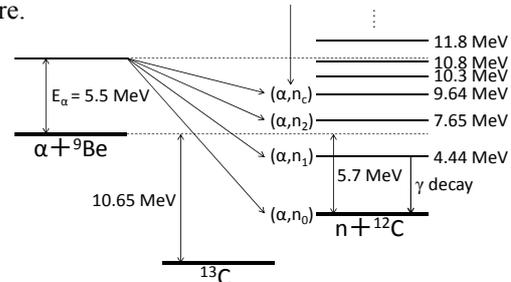


Figure 4. Energy level scheme of ^{12}C .

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