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**A new spent fuel source characterization code CHARS
and its application to the shielding of the thorium cycle**

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Radioprotection measures needed in the nuclear fuel cycle require accurate knowledge of the radioactive sources involved. For innovative nuclear reactors such as Generation IV designs radiation sources (alpha, beta, gamma and neutron) in the spent fuel need to be calculated in order to understand the radioprotection needed in all aspects of the fuel cycle (transport, reprocessing, fuel fabrication, waste storage, etc.). For this purpose we have developed CHARS a set of source characterization tools coupled to our code MURE (MCNP Utilities for Reactors Evolution) which is a precision research code for fuel evolution. MURE determines inventories of around 800 nuclei during irradiation and cooling via a series of MCNP5 calculations and numerical integration of Bateman's equations. With the CHARS package using the ENSDF libraries, it's possible to generate alpha, beta and gamma spectra of the nuclear fuel at the end of cycle from any given reactor design. These complex source definitions are then used to generate automatic MCNP5 inputs for radioprotection calculations, allowing us to undertake a wide range of radioprotection studies. The first use of these tools was to estimate additional shielding in the French fuel cycle in case of switching from the current uranium (U/Pu) cycle to the thorium (Th/U) cycle. Irradiated thorium-based fuels produce small quantities of ^{232}U , which has a relatively short half-life (69 years) and emits a hard gamma of 2.6 MeV at the end of its decay chain. ^{232}U is synthesized in mainly two ways: $^{233}\text{U}(n,2n)$ and $^{232}\text{Th}(n,2n)$ followed by $^{231}\text{Pa}(n,\gamma)$. From the results of radioprotection calculations we estimate additional thickness of shielding required for the back end of the fuel cycle and show that the greatest constraints occur for the fuel manufacturing. On the other hand the neutron yield for some thorium-based fuels will be lower in Th/U than in U/Pu cycle and as a result necessary neutronic shielding will be reduced.

Keywords: *spent fuel; thorium cycle; gamma spectra; neutron spectra; Monte Carlo; MCNP; shielding; dose rate; MURE*

1. Introduction

The next generation of nuclear reactors (Generation IV) is currently being developed and is intended to improve fissile resource utilization, reduce waste production, increase safety, and enhance anti-proliferation. The future use of Generation IV technologies requires detailed designs, not just of the reactors themselves, but their entire nuclear fuel cycle. Correct dimensions for the factories of the back end of the fuel cycle (fuel reprocessing, re-fabrication, transportation and waste disposal) need to be calculated along with the design of the radioprotection shielding needed. This requires accurate knowledge of the radioactive sources (alpha, beta, gamma and neutron) involved at each stage in the fuel cycle. In order to facilitate such calculations, we developed a code for source characterization (CHARS) coupled with the existing precise fuel depletion code [1] (MURE), which could produce accurate spent fuel

composition information for complex reactor geometries using different fuel cycles.

2. Tools used and developed

2.1. MURE

MURE is an interactive precision research code coupled with widely used Monte-Carlo based neutron, photon, electron, or coupled neutron/photon/electron transport code, MCNP [2]. MURE determines inventories of around 800 nuclei during fuel irradiation and cooling via a series of MCNP5 calculations and numerical integration of the set of coupled differential equations that govern their production and destruction (Bateman's equations). MURE is also a C++ interface for MCNP5 facilitating automated input/output, the geometry definition and post-processing of results. A graphical user interface (MureGui) allows access to the data generated during the fuel evolution calculation: k_{eff} , fluxes, decay heats, inventories, cross-sections, radio toxicities, reaction

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rates. In addition the MureGui now computes radiation spectra using the CHARS package.

2.2. CHARS (CHAracterization of Radioactive Sources)

CHARS uses the results of MURE's precision fuel depletion calculation and combines them with the best available nuclear decay data to reconstruct neutron, gamma, alpha and beta spectra of any spent fuel after irradiation. This includes during the cooling phase in-core, and before and after reprocessing out of core.

2.2.1 Gamma spectra

The ability to read nuclear structure data in ENSDF [3] format and the knowledge of nuclei inventories during cooling allows reconstruction of gamma spectra from the decays of all nuclei combined together. Gammas taken into account are gammas from α , β^- , β^+ , electron capture and isomeric transition decay processes. X rays originating from bremsstrahlung of β^- can be estimated using MCNP and beta spectra calculated using a simple Fermi theory where all the transitions are allowed (the details of this calculation are not discussed here since it is not important for fabrication of new fuel). The nuclear structure data used comes from [4] but it is possible to use any nuclear structure data library in the ENSDF format.

2.2.2 Neutron spectra

Neutrons emitted from the spent fuel are produced by two major processes, one is spontaneous fission and another is (α, n) reactions on light nuclei. In the CHARS code the neutron spectra from spontaneous fission is computed using a Watt distribution $W(E)$.

$$W(E) = C \exp\left(-\frac{E}{a}\right) \cdot \sinh(\sqrt{bE}) \quad (1)$$

Where a and b are coefficients depending on the nucleus and C the normalization constant. The neutron yield of the nucleus by spontaneous fission at energy E is:

$$A(E) = \lambda N \langle \nu \rangle BR \cdot W(E) \quad (2)$$

Where BR is the branching ratio of the spontaneous fission and $\langle \nu \rangle$ the average number of neutron per fission. Values of BR are from ENSDF, $\langle \nu \rangle$ from ENDF 7.1 and a, b are from a fit performed by Madland (see [5]). λ is the decay constant and N is the number of nuclei. Neutrons are also produced by (α, n) reactions on light nuclei during slowing down of alpha particle in the spent fuel. Two principal reactions $^{17}\text{O}(\alpha, n)^{20}\text{Ne}$ and $^{18}\text{O}(\alpha, n)^{21}\text{Ne}$ are taken into account for calculation of neutron yield attributed from (α, n) reactions. To compute neutron spectra from these reactions, one needs alpha spectra (calculated using ENSDF data [4]), stopping power in the media, cross sections leading to different level of ^{20}Ne and ^{21}Ne and oxygen density. The stopping powers are calculated for UO_2 media using the SRIM (Stopping and Range of Ions in Matter) [6] code and the total (α, n) cross-section (σ_{tot}) is taken from JENDL/AN-2005 [7]. The ratio $\frac{\sigma_i(E)}{\sigma_{tot}(E)}$, where σ_i is the

(α, n) cross section of the reaction leading to the level l of the daughter nucleus, is calculated with the nuclear reaction simulation code TALYS [8]. Thick target, homogenous mixture and isotropic emission of neutrons in the center of mass are assumed for the calculation.

2.3. Validation

2.3.1 Benchmark with CESAR 5.33

CESAR is a qualified depletion code developed by the CEA and AREVA (COGEMA) [9]. The benchmark consists of calculating the gamma spectrum of a standard pressurized water reactor (PWR) 17x17 uranium oxide (UOX) assembly, enriched to 5% mol. ^{235}U at a fuel burn-up of 55GWd/tHM and various different fuel cooling times.

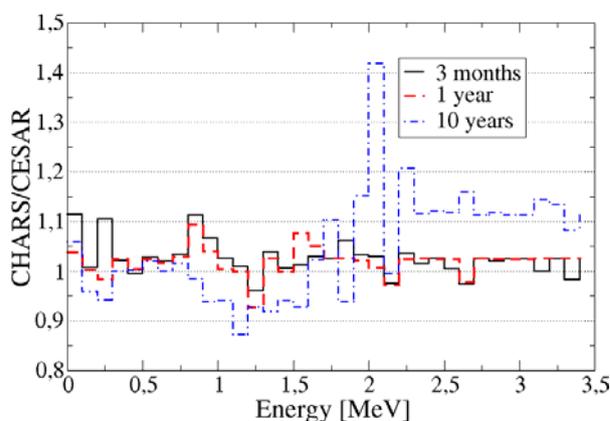


Figure 1. Ratio of spectra calculated by CESAR and CHARS for several cooling time.

Figure 1 shows good agreement between spectra calculated by these two codes for 3 months and 1 year of cooling. The two codes use the same database for the neutron transport part and fission yield (JEFF3.1) but use different ones for nuclear structure data: ENSDF for CHARS (except for ^{106}Rh which came from JEFF3.1) and JEFF3.1 for CESAR. The agreement is very good regarding the total gamma emission on a wide range of cooling time (**Table 1**). The discrepancy at 10 years for energies above 2 MeV comes from a greater production of ^{106}Rh with MURE and the contribution of this nucleus in the spectrum above 2MeV increases with time. In addition CESAR does not consider the nuclei ^{102}Rh and ^{208}Tl . This leads to the peaks in energy range 2-2.1 MeV, 2.2-2.3 MeV and 2.6-2.7 MeV.

Table 1. Gamma emission of a UOx spent fuel.

Cooling time	CESAR ($\gamma \cdot s^{-1} \cdot tHM^{-1}$)	CHARS ($\gamma \cdot s^{-1} \cdot tHM^{-1}$)
3 months	1.11×10^{17}	1.14×10^{17}
1 year	3.44×10^{16}	3.52×10^{16}
10 years	5.59×10^{15}	5.58×10^{15}
100 years	6.21×10^{14}	6.30×10^{14}

2.3.2 Neutron Spectra of PuO₂

In the work of E.F. Shores [10] a benchmark of the SOURCES code was performed using neutron spectra measurements of several plutonium oxides of various compositions. This benchmark was also used to validate CHARS. Results of neutron yields are presented in **Table 2**. References about measurements are available in [10].

Table 2. Measured and calculated neutron source magnitudes.

Sample	Measured [n.s ⁻¹]	CHARS [n.s ⁻¹]	Difference [%]
JAERI1	184.7±4.8 sf	197	6.41
	157.8±9.8 (α,n)	125	-20.9
	342.5±8.6 total	321	-6.19
JAERI2	42.1±1.1 sf	42	-0.19
	51.3± (α,n)	47	-7.60
	93.4±/-2.3 total	89	-4.66
ORNL	8.6785 x 10 ⁴ ± 867	7.63 x 10 ⁴	-12.0
SRS	8.55 x 10 ⁵ ±1.28 x 10 ⁴	8.62 x 10 ⁵	0.87
PNL	2.3 x 10 ⁵	2.35 x 10 ⁵	2.01

Table 2 shows an acceptable agreement between CHARS and the measurements with a discrepancy between -12% and 6.41% for the total yield. These differences are mainly due to the (α,n) part where the discrepancy can be high (-20.9% for the JAERI1 sample). Shapes of the spectra are in good agreement with those calculated by SOURCES as shown in **Figure 2**.

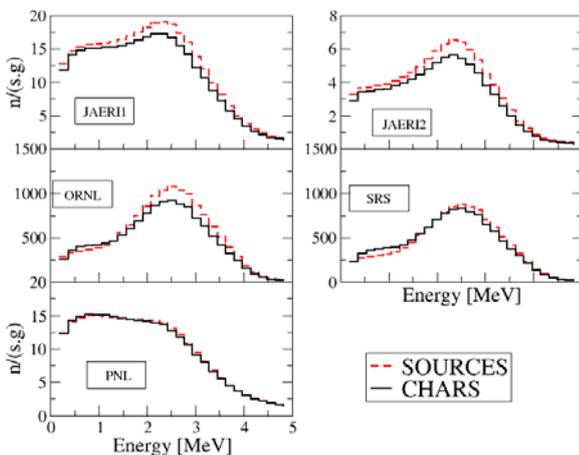


Figure 2. Neutron spectra calculated by CHARS and SOURCES for one gram of samples of plutonium dioxide.

3. Application

In order to save the fissionable natural resources, the use of the thorium cycle is a good candidate for replacing the current U/Pu cycle since the fissile nucleus of the cycle, ²³³U, has neutronic properties favorable to a much better regeneration of fissile material in thermal reactors. Moreover, as the nuclei of the thorium cycle are lighter than in the U/Pu cycle, the production of minor actinides is significantly reduced and hence waste inventories could

be much lower. However, the use of Th is only viable if the spent fuel is reprocessed to recover the fissile ²³³U, which does not exist in nature. This reprocessing will involve a heavy industrial infrastructure, particularly since thorium based spent fuel contains small quantities of ²³²U, which is the mother of the hard gamma emitter (²⁰⁸Tl) of 2.6 MeV. This gamma activity will certainly be a problem for the shielding of the back end fuel cycle especially for the manufacturing of new Th/U mixed oxide (MOX) fuel since the ²³²U (half life of 69 years) is present in the spent fuel uranium vector.

At present MOX fabrication plants such as MELOX in France use the MIMAS process for Pu/U based MOX fuels. This consists of two stages of blending. In the first one, called the primary blend, PuO₂ powder and UO₂ powder are blended together. At this point the amount of PuO₂ is 30 weight %. In the second stage, UO₂ is added to the mix in order to achieve the desired plutonium content. The secondary blend is then pelletized and sintered before being inserted into fuel rods and then fuel assemblies. The goal of this study is to show the impact on radiation exposure of the fabrication of ThU MOX fuel comparatively to the conventional UPu MOX fuel.

3.1. Problem definition

Due to its high radiation level the first stage of blending has been chosen for the comparison. The plutonium used for the MOX came from a PWR 5% ²³⁵U at 55GWd/tHM with 5 years of cooling before reprocessing and 6 years of cooling after. This elapsed time corresponds to an ²⁴¹Am content of 3%. The uranium for the ThU fuel came from a PWR loaded with ThPuO₂, 11% of Pu at 55 GWd/tHM with 5 years of cooling before reprocessing; several different cooling times after reprocessing are used in order to see the impact of the cooling time on the dose rate.

The design of the first blending workshop consists of a stainless steel glove box of thickness 1 cm with a lead glass window of 2 cm in thickness. The mixture, which is 60 kg of MOX or ThUO₂ is in a stainless steel bottle of 2 mm in thickness surrounded by a neutron protection layer of 10 cm thick polyethylene. The dose rate is estimated at 10 cm from the surface of glove-box window using MCNP5 and the track-length estimator (F4 tally) modified by the AP fluence-to-dose conversion factor from ICRP [11].

3.2. Source calculations

Given the composition of the first blend, CHARS was used to calculate both gamma and neutron spectra. The composition assumed two oxygen atoms for each heavy nucleus. The thorium used for the ThUO₂ powder of 30 weight % of UO₂ was pure ²³²Th, and depleted uranium (2.01 %mol. ²³⁵U) was used for the MOX powder containing 30 weight % of PuO₂. Results of this calculation are shown in **Figure 3** for gamma spectra. The origin of gammas and neutrons for the ThUO₂ powder is the ²³²U decay chain while for the MOX case is mainly ²³⁸Pu and ²⁴⁰Pu for neutrons and ²⁴¹Am for gammas. Total

neutron emission is lower in the ThUO₂ case but the spectrum is slightly harder due to the highest energy of alphas (up to 8.78 MeV (from ²¹²Po)) whereas the maximum energy for the UPuO₂ standard MOX case is 5.5 MeV (from ²³⁸Pu)).

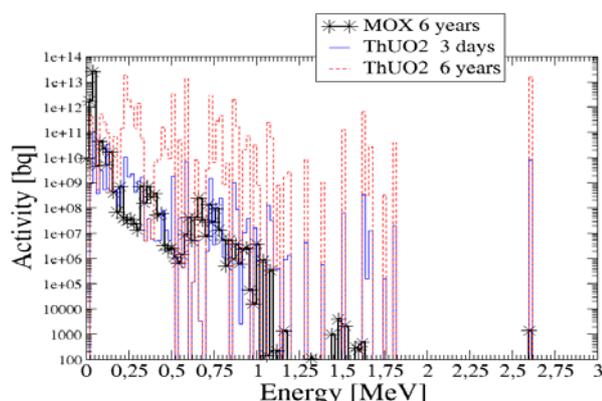


Figure 3. Gamma Spectra of the two different powders.

3.3. Results

Table 3 clearly shows the problem caused by the ²³²U in the fabrication of ThU fuels. It presents the dose rate behind the glove-box when using ThUO₂ powder of several cooling times relative to the dose rate calculated for the reference MOX powder after 6 years of cooling; i.e ThUO₂ dose rate / MOX dose rate. It appears that even if ThUO₂ fuel is made 3 days after reprocessing the dose rate is still 8 times higher than in the reference case. In the worst case: when the quantity of ²⁰⁸Tl is at its maximum, the dose rate is 16500 times higher. On the other hand, the neutron dose rate is reduced and a lighter neutron shielding can be used.

Table 3. Relative dose rate (ThUO₂ / MOX).

Time after reprocessing	Gamma (Secondary γ)	Neutron	Total
3 days	149 (0.083)	0.076	8
1 months	9990 (0.103)	0.089	536
1 year	1.02x10 ⁵ (0.237)	0.208	5470
6 years	2.9x10 ⁵ (0.51)	0.447	15600
10 years	3.08x10 ⁵ (0.53)	0.467	16500

In the next calculation, the cooling time is kept to 10 years (worst case) and the lead glass thickness is increased. The results are presented in **Table 4**. In order to obtain the same dose rate as the reference case the lead glass window must be 43.5 cm, which is not realistic.

Table 4. Relative dose rate versus lead glass thickness.

Thickness (cm)	2	5	10	20	30	43.5
Relative dose rate	16500	8330	2700	266	25	1

4. Conclusion

A set of tools to characterize spent fuel has been developed and its reliability has been shown with several benchmarks. CHARS will be available at the OECD NEA at the end of 2013 at the time of the update of the MURE package. With this tool and MCNP, we demonstrate in a qualitative way that the fabrication of ThUO₂ fuel is not possible on an industrial scale using currently used processes such as MIMAS in glove-boxes. Complete automation of the fuel re-fabrication process may be required which could place severe problems for long-term maintenance of the facility. Research into other fabrication processes is clearly needed (e.g. impregnation of low density ThO₂ pellets with uranyl nitrate [12]).

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