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Shielding Experiments at High Energy Accelerators of Fermilab (III): Neutron Spectrum Measurements in Intense Pulsed Neutron Fields of The 120-GeV Proton Facility Using A Current Bonner Sphere Technique

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Neutron energy spectra behind iron and concrete shields have been measured at the anti-proton production target station in Fermi National Accelerator Laboratory to provide shielding benchmark data of neutrons produced from 120-GeV proton induced reaction. To measure neutron spectra in the radiation bursts accompanied by the injection of intense pulsed proton beam into the target, we have developed a current readout technique of a multi-moderator spectrometer. We employed a pair of BF₃ proportional counters filled with different ¹⁰B isotope enrichment gases, ^{nat}BF₃ (18% enrichment) and ¹⁰BF₃ (96% enrichment) and used them as current readout ionization chambers to avoid signal pile-up problem. The spectrometer is aimed to discriminate neutrons from other particles such as γ -rays and muons incoming in short bursts by subtracting results of the ^{nat}BF₃ counter from that of the ¹⁰BF₃ counter. The neutron spectrum obtained with the present technique is generally consistent with results of neutron-induced activation measurements.

KEYWORDS: neutron energy spectra, multi-moderator spectrometer, radiation bursts, shielding benchmark data

I. Introduction

Estimation of neutron production and transmission through shielding are very important at high energy accelerator facilities because neutrons have high penetrability through shielding and large contribution to the external dose. A number of shielding design of high energy accelerator facilities are performed using particle interaction and transport codes such as MCNPX, FLUKA, MARS, PHITS and so on.¹⁻⁴⁾ These codes should be assessed based on benchmark measurements to ensure the calculation results. However, benchmark data on neutron energy spectra behind shields induced by high energy protons above 1 GeV are few because of lack of suitable facilities with a simple shielding structure and difficulties of neutron spectrum measurements especially for radiation bursts accompanied by beam operation with short pulse.⁵⁻⁹⁾

We have performed series of shielding experiments at the anti-proton production target facility (pbar target station) in FNAL to measure the induced activities and neutron flux around high energy accelerators in the Japanese and American Study of Muon Interaction and Neutron detection (JASMIN) collaboration.¹⁰⁻¹²⁾ The pbar target station has developed to provide anti-protons to the proton - anti-proton collider (TEVATRON) using 120 GeV protons hitting the

target. This facility has a simple structure above the target and beam dump and source term can be clearly defined. On the other hand, intense pulsed proton beam impacted on the target with power of about 65 kW and pulse width of several micro seconds produce radiation bursts even above thick shields, which make it difficult to apply active counters to measurements of the neutron spectrum.

In this study, we have developed a current readout Bonner sphere technique to measure neutron spectra in radiation bursts at the pbar target station by using a multi-moderator spectrometer (Bonner spheres) using a pair of BF₃ proportional counters filled with different ¹⁰B isotope enrichment gases, ^{nat}BF₃ (18% enrichment) and ¹⁰BF₃ (96% enrichment). The combination of different ¹⁰B-enriched BF₃ counters allow to discriminate the signals of neutrons from other particles such as γ -rays and muons by subtracting results of the ^{nat}BF₃ counter from those of the ¹⁰BF₃ counter. The neutron energy spectra behind shields are deduced by a conventional unfolding technique. The measured neutron energy spectra were compared with the results of neutron-induced activation measurements by means of reaction rates of nuclides.¹¹⁾

II. Method

1. Detector

Figure 1 shows a schematic view of the spectrometer. The spectrometer consists of two 2-inch-diameter BF₃

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spherical proportional counters (LND 2708) and four spherical polyethylene moderators with different diameters. One of these BF₃ counters is filled with 96% ¹⁰B-enriched BF₃ gas (¹⁰BF₃) and the other is 18% ¹⁰B-enriched BF₃ gas (^{nat}BF₃). The gas pressure of these BF₃ counters is 760 Torr. The difference of outputs from ¹⁰BF₃ counter and ^{nat}BF₃ counter is then considered to be only due to neutrons, since ¹⁰B has a much higher capture cross-section to low-energy neutrons than that of the other stable isotope ¹¹B.

The spherical moderators are made of polyethylene (density of 0.928 g/cm³) and their diameters are 81, 110, 150, 230 mm. These moderators were originally designed by Uwamino et al. to use with a 5 atm ³He proportional counter. ¹³ The ³He counter was replaced with ¹⁰BF₃ counter for this study to increase output charge produced in the neutron reaction owing to the higher Q-value of the ¹⁰B(n,α)⁷Li reaction (2.79 MeV) than that of the ³He(n,p)³H reaction (0.764MeV), which is important to use it as a current readout ionization chamber described in later.

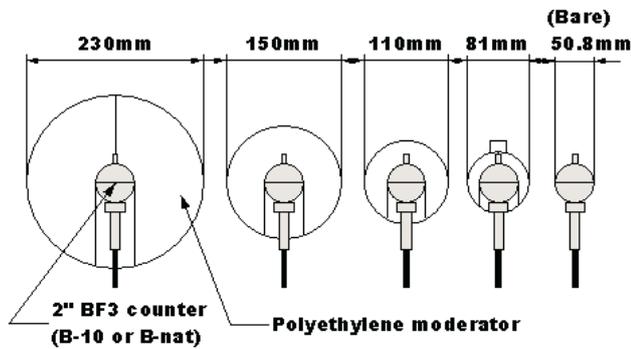


Fig. 1 The multi-moderator spectrometer used in this experiment. The spectrometer system has ten setups consist of two BF₃ counters (¹⁰BF₃ and ^{nat}BF₃) and five moderator patterns (0, 81, 110, 150, 230 mm)

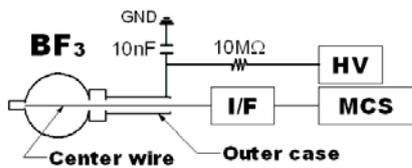


Fig. 2 Electronics of the current readout mode for a BF₃ counter. High voltage is applied to the outer case of the BF₃ counter to operate it as a current readout ionization chamber

2. Signal readout

Figure 2 shows a schematic diagram of electronics to operate the BF₃ counter as a current readout ionization chamber. In the current readout mode, high voltage was applied to the outer case of the BF₃ counter to read the current induced by ¹⁰B(n,α)⁷Li reaction on the central wire. The applied voltage was determined to be 300 V, based on a measured saturation curve of the induced charge and leakage current. The current was fed to a current to frequency converter (I/F; Laboratory Equipment Co.), which converts input charge of 1 pC to 1 pulse, and recorded by a

multi-channel scaler (MCS) to measure the signals synchronized with the beam injection and monitor the fluctuation.

3. Response functions

Figure 3 shows the response functions for ¹⁰BF₃ counters with polyethylene moderators of 0, 81, 110, 150, 230 mm diameters calculated using the MCNPX code (version 25f) ¹¹ with the evaluated cross-section data libraries. The library used for neutron transport and reaction rate were LA150 (hydrogen, carbon and stainless steel) and ENDF/B-VI (boron and fluorine) ¹⁴. For the energy region that is not covered by these libraries, MCNPX default physics models were used. The calculation was performed with mono-energetic parallel neutron beam of different energies from 10⁻⁹ eV to 400 MeV for each moderator diameter.

The response functions were measured with a conventional pulse readout system consists of a pre-amplifier and an amplifier and a multi-channel analyzer using 8.0 keV monoenergetic neutrons at the FRS facility in Japan Atomic Energy Agency. ¹⁵ The agreement between experimental and calculated values was rather good as shown in **Fig. 3**. The shape of response functions in the current readout mode should not be so different with that in the pulse readout mode, because the both output signals are based on the measurement of the charge ionized by the outgoing particles from the ¹⁰B(n,α)⁷Li reaction. Iijima et al. report that the shape of the response function in a current readout mode agrees well with the calculation results of MCNPX as well as the response function in a pulse readout mode using a ³He counter covered with polyethylene moderator. ¹⁶

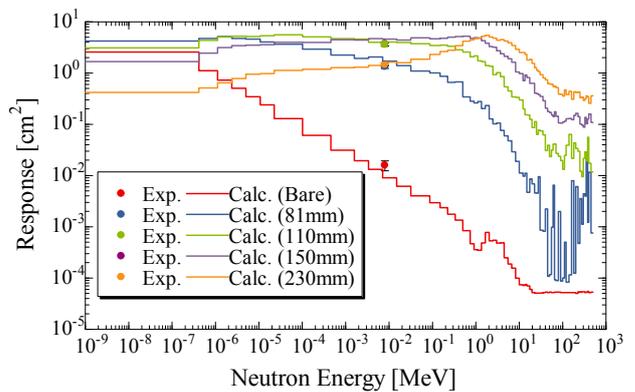


Fig. 3 Response functions of the ¹⁰BF₃ counter with five-moderators for neutron energies from 10⁻³ eV to 400 MeV

4. Calibration

To obtain the conversion factor from the current output (pC) to corresponding number of the ¹⁰B(n,α)⁷Li reaction in the BF₃ counters, we irradiate a pair of ¹⁰BF₃ counters covered with the 230-mm-thick moderator using a neutron reference source of 3-Ci ²⁴¹Am-Be and measure the output counts per incident neutron in the pulse readout mode and the current readout mode. The ²⁴¹Am-Be neutron source was placed 74 cm and 17 cm away from the center of the spectrometer in the measurement of the pulse readout mode

and the current readout mode, respectively.

The events induced by neutrons and γ -rays are clearly distinguished from their pulse height distribution. The integrated charge of γ -rays is estimated to be less than 0.1 % of that of neutrons owing to the high Q-value of the $^{10}\text{B}(n,\alpha)^7\text{Li}$ reaction. Although the contamination of γ -rays to the neutron measurement in the current readout mode is small, the source were covered with a 2-mm-thick lead cap in the measurement of conversion factor to suppress the γ -ray contamination. The conversion factors which were obtained from ratio of count rates between pulse mode and current mode is 56.6 and 61.8 pulse/pC for the $^{10}\text{BF}_3$ counter and the $^{nat}\text{BF}_3$ counter, respectively.

III. Experiment

Figure 4 shows geometry of the pbar target station shielding structure. The station consists of anti-proton production target made of multiple layers of Inconel and copper disks and graphite dump with iron and concrete shielding. The simple structure and material of the station is suitable to provide benchmark data of neutron penetration. The anti-proton production target is irradiated with 120 GeV proton beam provided from the Main Injector ring. The toroidal coil (TR105B) is placed just before the antiproton production target to monitor the beam intensity during this experiment. The typical time structure of the proton beam is 1.6 μs pulse width, with a cycle time of 3.0 s. The power of the beam reached as high as 65 kW during our experiment. Downstream of the target, a graphite/aluminum dump is placed to absorb the remaining protons and secondary particles after the some extraction devices of produced anti-protons into the pbar extraction channel. These components are surrounded by 183-cm-thick iron and 122-cm-thick concrete shields. The area above these components is a restricted area during beam operation. The spectrometers of the ^{10}B counter and the ^{nat}B counter were placed at two points that is 50 cm aside from the beam axis above the target shown in **Fig.4**. Activation samples of In, Al, Cu, Bi were placed at the point on the beam axis above the target for the other measurement of neutron-induced activities after 10 days irradiation.¹¹⁾

IV. Results

Figure 5 shows a result of the MCS output of the $^{10}\text{BF}_3$ counter with the 110-mm-diameter moderator located above the target. The measured time profile of the current induced in the counter was clearly synthesized with each injection of pulsed beam obtained from TR105B. The mean current was obtained by integrating the pulse data which are subtracted with the background current which is ~ 5 pA due to detector leakage current. The neutron spectra were deduced using an unfolding code, SANDII¹⁷⁾, with the response function and an initial guess spectrum calculated by the MARS15 code.¹²⁾ The mean count rates of the $^{10}\text{B}(n,\alpha)^7\text{Li}$ reaction in the counter were obtained by subtracting results of the $^{nat}\text{BF}_3$ counter from that of the $^{10}\text{BF}_3$ counter.

Figures 6 show the neutron energy spectra obtained together with the spectrum calculated by the MARS15 used

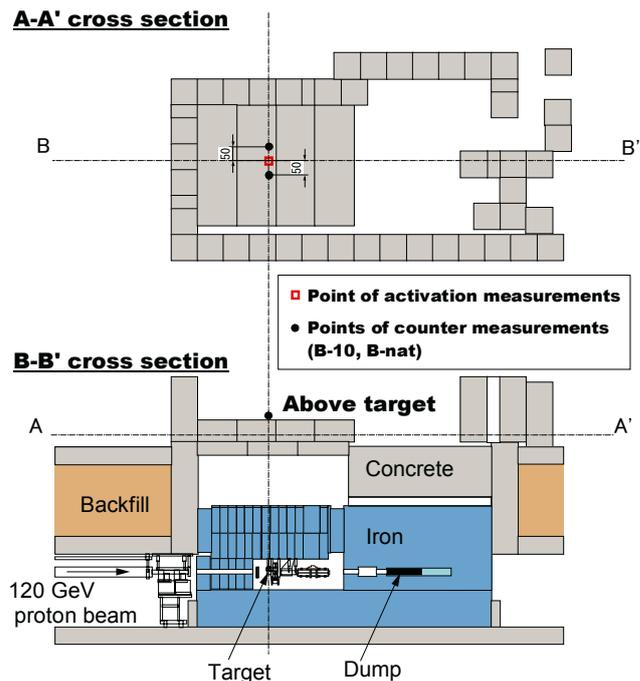


Fig. 4 Geometry of pbar target station

for the initial guess spectrum in the unfolding analysis. The

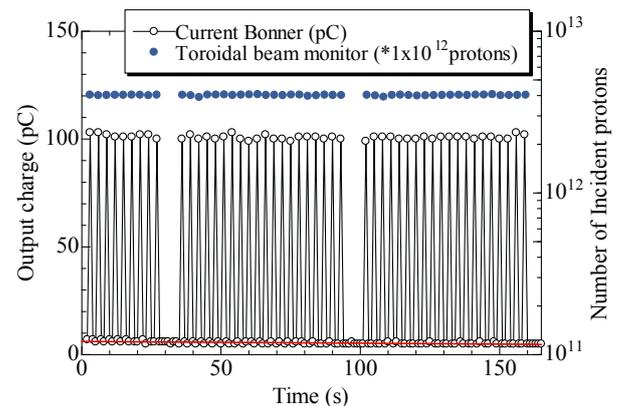


Fig. 5 MCS output obtained from $^{10}\text{BF}_3$ counter covered with 110-mm-thick moderator located above the target. The beam data of TR105B are also plotted.

measured spectrum has errors of the TR105R beam monitor (10%) and the I/F converter (10%), and the statistical errors (9.3%). The total error of the neutron flux was estimated to be 17.1 % without ambiguity of the unfolding method. The calculation generally agrees with the experimental results, but underestimates thermal neutron flux and slightly overestimates fast neutron flux above 1 MeV. The difference of thermal neutron flux which is drastically affected by surroundings and gaps in shielding would be due to ambiguities on the geometry description in the MARS calculation. The difference of fast neutron flux would be due to ambiguities of densities and compositions of materials described in the MARS calculation.

Figure 7 shows the comparison of reaction rates of nuclide between two experiments. Present results were deduced from the measured neutron spectra multiplying the corresponding

activation cross-sections shown in ref.11. The reaction thresholds of $^{115}\text{In}(n,n')$, $^{27}\text{Al}(n,\alpha)$, $^{209}\text{Bi}(n,4n)$, $^{209}\text{Bi}(n,5n)$, $^{209}\text{Bi}(n,6n)$ are 0.6, 3.3, 22.6, 29.6, 38, 45.3 MeV, respectively. Present results show in good agreement with the activation results within 30 % in the energy range.

V. Conclusion

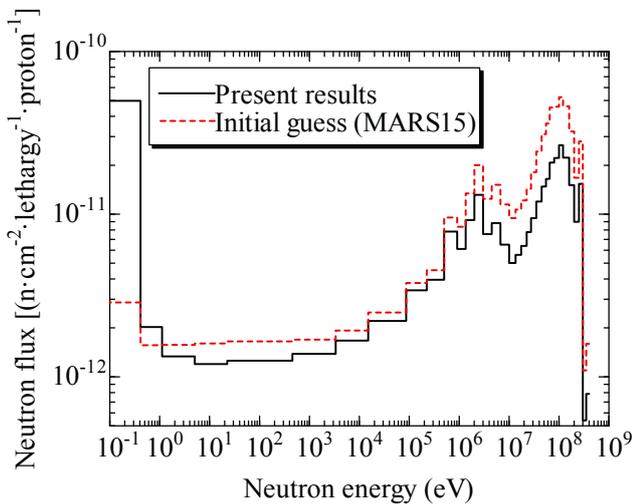


Fig. 6 Neutron energy spectra above the target unfolded with the initial spectrum calculated by the MARS code

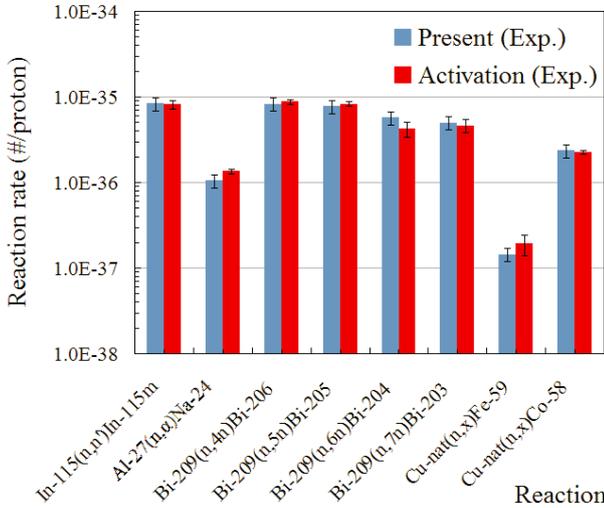


Fig. 7 Comparison of reaction rates of nuclides between present results and activation results

We have developed a current readout Bonner sphere technique to measure neutron spectra in radiation bursts at the pbar target station by a multi-moderator spectrometer (Bonner spheres) using a pair of BF_3 proportional counters filled with different ^{10}B isotope enrichment gases, $^{nat}\text{BF}_3$ (18% enrichment) and $^{10}\text{BF}_3$ (96% enrichment) and measured the neutron energy spectra behind 183-cm-thick iron and 122-cm-thick concrete shields from 120 GeV proton induced reactions. Present result was consistent with the other experimental results by the activation method which shows reliability of both data. The initial guess spectrum

calculated by MARS15 show slightly overestimation in MeV region and underestimation in the thermal energy region, which would be due to ambiguities on the geometry description in the MARS calculation should be investigated further.

The current readout Bonner sphere technique would be helpful to measure shielding benchmark data in high energy proton accelerators under the intense burst field that is impossible to use active counter in pulse mode. Although the lack of sensitivity for over 10 MeV neutrons should be improved in further experiment, the spectrum results would be helpful to investigate particle interaction and transport codes for high energy region.

Acknowledgment

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